

# Short Takes 331

Green's functions  
of the Laplacian:  
eigenfunction  
expansion

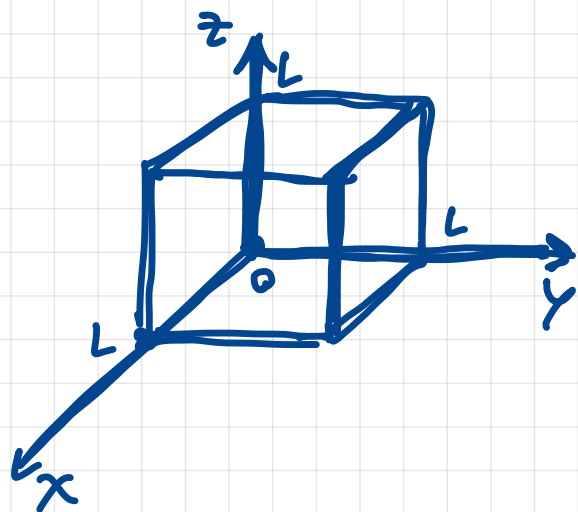


# Green's functions of the Laplacian: eigenfunction expansion.

$$\nabla^2 G(\vec{r}, \vec{r}') = \delta(\vec{r} - \vec{r}') \quad \& \quad \text{b.c}$$

In the last few videos we used  $\nabla^2 \phi_{\vec{n}} = \lambda_{\vec{n}} \phi_{\vec{n}}$  for...

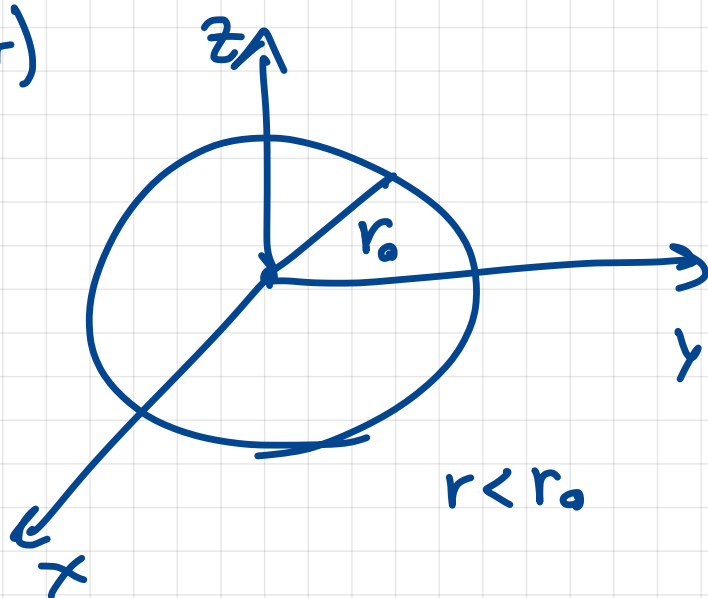
- Cube  
(Cartesian coords)



- $f(\vec{r}, t) = 0$  on faces of the cube.

$$\left[ \begin{array}{l} \phi_{\vec{n}}(\vec{r}) = \phi_{n_1}(x) \phi_{n_2}(y) \phi_{n_3}(z), \quad \vec{n} = (n_1, n_2, n_3), \quad \phi_{n_i}(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n_i \pi x}{L}\right) \\ \lambda_{\vec{n}} = \lambda_{n_1} + \lambda_{n_2} + \lambda_{n_3}, \quad n_i = 1, 2, 3, \dots, \quad \lambda_{n_i} = -\left(\frac{n_i \pi}{L}\right)^2 \end{array} \right.$$

- Sphere (interior)  
(spherical coords)



- $f(\vec{r}, t)|_{|\vec{r}|=r_0} = 0$

$$\left[ \begin{array}{l} \phi_{\vec{n}}(\vec{r}) = R_{\lambda_n}(r) Y_{lm}(\theta, \varphi), \quad \vec{n} = (n, l, m) \\ \lambda_{\vec{n}} = -\left(\frac{\chi_{ln}}{r_0}\right)^2 \end{array} \right. \quad \begin{array}{l} n = 1, \dots \\ l = 0, 1, \dots \\ m = -l, \dots, l \end{array}$$

$$R_{\lambda_n}(r) = j_l\left(\chi_{ln} \frac{r}{r_0}\right), \quad \chi_{ln}: n\text{-th (non-zero) root of } j_l(x).$$

"spherical Bessel"

As we've seen in previous videos,

$$G(\vec{r}, \vec{r}') = \sum_{\vec{n}} \frac{\phi_{\vec{n}}^*(\vec{r}') \phi_{\vec{n}}(\vec{r})}{\lambda_{\vec{n}}},$$

since

$$\begin{aligned} \nabla_{\vec{r}}^2 G(\vec{r}, \vec{r}') &= \sum_{\vec{n}} \frac{\phi_{\vec{n}}^*(\vec{r}') \nabla_{\vec{r}}^2 \phi_{\vec{n}}(\vec{r})}{\lambda_{\vec{n}}} \\ &= \sum_{\vec{n}} \phi_{\vec{n}}^*(\vec{r}') \phi_{\vec{n}}(\vec{r}) = \delta(\vec{r} - \vec{r}') \quad \checkmark \end{aligned}$$

"completeness relation"

cube

$$\underbrace{\left[ \sqrt{\frac{2}{L}} \sum_{n_1} \sin\left(\frac{n_1 \pi x}{L}\right) \sin\left(\frac{n_1 \pi x'}{L}\right) \right]}_{\delta(x-x')} \underbrace{\left[ \sqrt{\frac{2}{L}} \sum_{n_2} \dots \right]}_{\delta(y-y')} \underbrace{\left[ \sqrt{\frac{2}{L}} \sum_{n_3} \dots \right]}_{\delta(z-z')}$$

Note  $G$  does not factorize, though!

sphere

$$\sum_{n, l, m} R_{ln}(r) Y_{lm}(\theta, \varphi) R_{ln}^*(r') Y_{lm}^*(\theta', \varphi') = \delta(\vec{r} - \vec{r}')$$

much less obvious!

What happens if we remove the boundaries?

Then,

$$\phi_{\vec{k}}(\vec{r}) \rightarrow \frac{1}{(2\pi)^{3/2}} e^{i\vec{k} \cdot \vec{r}} = \phi_{\vec{k}}(\vec{r})$$

$\lambda_{\vec{n}} \rightarrow -\hbar^2 = \lambda_{\vec{k}}$   $\vec{k}$  is now continuous!  
 (" $k_x, k_y, k_z$ ")

and

$$G(\vec{r}, \vec{r}') = \int d^3k \frac{\phi_{\vec{k}}(\vec{r}) \phi_{\vec{k}}^*(\vec{r}')}{\lambda_{\vec{k}}}$$

$$= \frac{1}{(2\pi)^3} \int d^3k \frac{e^{i\vec{k}\cdot(\vec{r}-\vec{r}')}}{-k^2}$$

$$\nabla^2 G(\vec{r}, \vec{r}') = \frac{1}{(2\pi)^3} \int d^3k e^{i\vec{k}\cdot(\vec{r}-\vec{r}')} = \delta(\vec{r}-\vec{r}')$$

→ Fourier transforms!

