

Short Takes

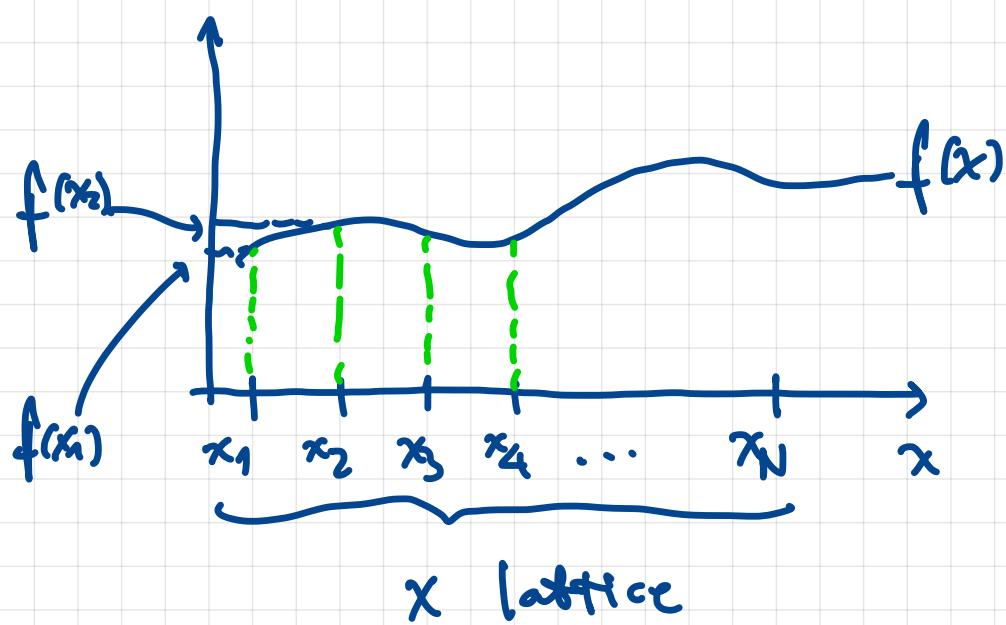
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Diagonalizing
differential
operators
on the lattice



Diagonalizing differential operators on the lattice

- How do we represent $f(x)$ on a computer?



$$f(x) \rightarrow \begin{pmatrix} f(x_1) \\ f(x_2) \\ \vdots \\ f(x_N) \end{pmatrix} = \vec{f}$$

or simply f_1, \dots, f_N

- How do we represent derivatives $f'(x), f''(x)$, etc?

$$\frac{df}{dx} = \lim_{h \rightarrow 0} \frac{f(x+h) - f(x)}{h}$$

on the lattice $\frac{f(x_{i+1}) - f(x_i)}{x_{i+1} - x_i}$

"forward derivative"

or, $\frac{f_{i+1} - f_i}{l}$

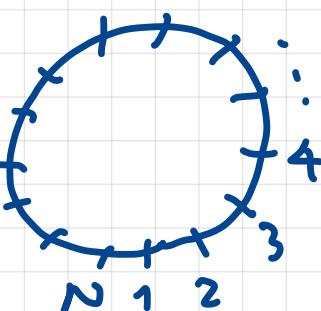
, assuming uniform lattice
where $x_{i+1} - x_i = l$

lattice spacing.

- As a linear operator, ie a matrix...

$$\frac{df}{dx} \rightarrow \frac{1}{l} \begin{pmatrix} -1 & 1 & 0 & \dots & 0 \\ 0 & -1 & 1 & \ddots & 0 \\ \vdots & 0 & \ddots & \ddots & \vdots \\ 0 & 0 & \dots & 0 & -1 \\ 1 & 0 & \dots & 0 & 1 \end{pmatrix} \begin{pmatrix} f_1 \\ f_2 \\ \vdots \\ f_N \end{pmatrix}$$

$N \times N$
rep of $\frac{d}{dx} \rightarrow D^{(1)}$



boundary conditions? → (periodic in this case)

Usually, when solving diff eqs, we need to think about these separately. Here, it comes with the matrix!

Notice...

- $D^{(1)} \begin{pmatrix} c \\ c \\ \vdots \\ c \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ \vdots \\ 0 \end{pmatrix}$, for any c
 $\rightarrow \text{Ker } D^{(1)}$ is non-trivial,
so we don't have a well-defined inverse op.
 Makes sense!

$$\frac{dc}{dx} = 0 \quad \forall c$$

- $[D^{(1)}]_{ij} = \frac{1}{\ell} (-\delta_{ij} + \delta_{i+1,j})$ (with periodic indexing
 $N+1 \leftrightarrow 1$)

$$\underbrace{D^{(1)} \vec{v}_h}_{\downarrow} = ? \quad , \quad [\vec{v}_h]_j = \frac{1}{\sqrt{N}} \exp\left(i \frac{2\pi}{N} h \cdot j\right)$$

$$[\]_q = \sum_{j=1}^N \frac{1}{\ell} (-\delta_{q,j} + \delta_{q+1,j}) \cdot \frac{1}{\sqrt{N}} \exp\left(i \frac{2\pi}{N} h \cdot j\right)$$

$$= \frac{1}{\ell} \frac{1}{\sqrt{N}} \left[-\exp\left(i \frac{2\pi}{N} h \cdot q\right) + \exp\left(i \frac{2\pi}{N} h \cdot (q+1)\right) \right]$$

$$= \frac{1}{\ell} \left[\exp\left(i \frac{2\pi}{N} h\right) - 1 \right] \frac{1}{\sqrt{N}} \exp\left(i \frac{2\pi}{N} h \cdot q\right)$$

$$= \lambda_h \cdot [\vec{v}_h]_q$$

$\rightarrow \boxed{\vec{v}_h \text{ is an eigenvector of } D^{(1)}}$

with eigenvalue $\lambda_h = \frac{1}{\ell} \left[\exp\left(i \frac{2\pi}{N} h\right) - 1 \right]$

"Fourier modes are eigenvectors of derivative ops."
 (check boundaries!)

- Note $\lambda_N = 0$ & $v_N = \frac{1}{\sqrt{N}}$ \rightarrow constant! (makes sense)

- Finally,

$$\vec{v}_{k'}^+ D^{(1)} \vec{v}_k = \lambda_k \delta_{kk'} \quad \text{because } \vec{v}_k^+ \vec{v}_k = \delta_{kk'}$$

Therefore

$$U = \begin{pmatrix} \vec{v}_1 & \vec{v}_2 & \dots & \vec{v}_N \end{pmatrix}$$

is the unitary matrix that
diagonalizes $D^{(1)}$

$$U^+ D^{(1)} U = D = \begin{pmatrix} \lambda_1 & & & \\ & \lambda_2 & & \\ & & \ddots & \\ & & & \lambda_N \end{pmatrix}$$

$$U_{kj} = \frac{1}{\sqrt{N}} e^{i \frac{2\pi}{N} k \cdot j}$$

our old Fourier matrix!

